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In Situ Oxygen Doping of Monolayer MoS₂ for Novel Electronics

Jian Tang, Zheng Wei, Qinqin Wang, Yu Wang, Bo Han, Xiaomei Li, Biying Huang, Mengzhou Liao, Jieying Liu, Na Li, Yanchong Zhao, Cheng Shen, Yutuo Guo, Xuedong Bai, Peng Gao, Wei Yang, Lan Chen, Kehui Wu, Rong Yang,* Dongxia Shi, and Guangyu Zhang*

In 2D semiconductors, doping offers an effective approach to modulate their optical and electronic properties. Here, an in situ doping of oxygen atoms in monolayer molybdenum disulfide (MoS₂) is reported during the chemical vapor deposition process. Oxygen concentrations up to 20–25% can be reliable achieved in these doped monolayers, $MoS_{2-x}O_x$. These oxygen dopants are in a form of substitution of sulfur atoms in the MoS_2 lattice and can reduce the bandgap of intrinsic MoS_2 without introducing in-gap states as confirmed by photoluminescence spectroscopy and scanning tunneling spectroscopy. Field effect transistors made of monolayer $MoS_{2-x}O_x$ show enhanced electrical performances, such as high field-effect mobility ($\approx 100 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$) and inverter gain, ultrahigh devices' on/off ratio (>10⁹) and small subthreshold swing value ($\approx 80 \text{ mV dec}^{-1}$). This in situ oxygen doping technique holds great promise on developing advanced electronics based on 2D semiconductors.

1. Introduction

2D transition metal dichalcogenides (TMDs) are emerging lowdimensional materials with potential applications for optics and electronics.^[1,2] Among many 2D-TMDs, 2H phase monolayer MoS_2 is a representative direct bandgap semiconductor with bandgap (E_g) of ~2.2 eV and has been intensely explored as an alternative for high-performance thin film transistors and largearea integrated electronics.^[3,4] In order to realize the full potential of 2D-MoS₂ in such applications, it requires to effectively modulate their electrical, optical, and structural properties by introducing impurities for doping.^[5-10] So far, the chemisorption and charge transfer through surface functionalization are effective approaches to achieve doping, but it has a relatively weak influences on band structures or improvements of electronic properties.^[11,12] By contrast, substitutional doping is more stable and capable of tailoring the bandgap of 2D-TMDs without introducing in-gap states.^[13-16] Such substitutional doping could be realized through replacing the host transition metal or chalcogen atoms with other elements during synthesis to form alloys such as $Mo_{1-x}W_xS_2$, $MoS_{2x}Se_{2(1-x)}$, and $V_x W_y Mo_{1-x-y} S_{2z} Se_{2(1-z)}$, etc.^[17–21]

Recently, oxygen doping of 2D-MoS₂ has attracted considerable research inter-

ests. Previous studies show that post treatments of intrinsic MoS₂ in air, ozone, or oxygen plasma could induce oxygen doping, evidenced from the enhanced catalytic reactivity, quenched photoluminescence, and improved electrical conductivity.^[9,18,22–24] Such doping processes usually lead to oxygen substitution and oxidation as well and consequently yield highly disordered or fragmented lattice structures. In situ oxygen substitutional doping in 2D-MoS₂ with a controlled and nondestructive manner is thus highly desired to preserve its original lattice structure, but still remains challenging so far.

J. Tang, Z. Wei, Q. Wang, Y. Wang, Dr. X. Li, B. Huang, Dr. M. Liao, J. Liu, Dr. N. Li, Y. Zhao, Dr. C. Shen, Y. Guo, Prof. X. Bai, Prof. W. Yang, Prof. L. Chen, Prof. K. Wu, Prof. R. Yang, Prof. D. Shi, Prof. G. Zhang Institute of Physics and Beijing National Laboratory for Condensed Matter Physics Chinese Academy of Sciences	J. Tang, Z. Wei, Q. Wang, Y. Wang, Dr. X. Li, B. Huang, J. Liu, Dr. N. Li, Y. Zhao, Dr. C. Shen, Y. Guo, Prof. X. Bai, Prof. W. Yang, Prof. L. Chen, Prof. K. Wu, Prof. D. Shi, Prof. G. Zhang School of Physical Sciences University of Chinese Academy of Sciences Beijing 100190. China	
Beijing 100190 China	Dr. B. Har Dr. V. Li Brof. D. Coo	
E-mail: ryang@iphy.ac.cn; gyzhang@iphy.ac.cn	Dr. N. Li, Prof. W. Yang, Prof. L. Chen, Prof. K. Wu, Prof. R. Yang, Prof. G. Zhang Songshan Lake Materials Laboratory	
The ORCID identification number(s) for the author(s) of this article can be found under https://doi.org/10.1002/smll.202004276.		
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after original online publication.	Prof. W. Yang, Prof. R. Yang, Prof. D. Shi, Prof. G. Zhang	
DOI: 10.1002/smll.202004276	Beijing Key Laboratory for Nanomaterials and Nanodevices Beijing 100190, China	

2. Results and Discussion

In this work, we demonstrated the in situ oxygen substitutional doping of monolayer MoS₂ in the chemical vapor deposition (CVD) growth process. By introducing oxygen gas in the growth environment, we can achieve uniform oxygen doping in monolayer MoS₂ with tunable dopant concentrations, i.e., $MoS_{2-x}O_{x}$, without changing the lattice ordering. We also show the creation of in-plane heterostructures of $MoS_{2-x}O_x$ with different doping levels by tuning the oxygen gas partial pressure during the growth process. The atomic and electronic structures of the $MoS_{2-x}O_x$ samples were systematically characterized by scanning tunneling microscopy (STM), scanning tunneling spectroscopy (STS), scanning transmission electron microscope (STEM), X-ray photoelectron spectroscopy (XPS), Raman spectroscopy, and photoluminescence (PL) spectroscopy. More importantly, we found that the bandgap of monolayer $MoS_{2-r}O_r$ can be tuned by the oxygen doping levels. Electrical measurements on the fabricated field effect transistors (FETs) and logic gates reveal that, at an optimized doping level (within 20% doping concentration), $MoS_{2-x}O_x$ monolayers feature excellent device performances including large current capacities, enhanced sheet conductance, and associated carrier mobilities. These superior electrical properties of monolayer $MoS_{2-x}O_x$ suggest their potentials for high performance and low power consumption electronics.

The in situ growth of $MoS_{2-x}O_x$ monolayers was performed in our home-made CVD system as illustrated in Figure 1a. MoO₃, S, and O_2 were used as reaction sources and the growth of $MoS_{2-x}O_x$ monolayers was achieved on preannealed c-plane sapphire substrates, which were vertically placed in the growth chamber. Such vertical placement is beneficial to achieve more uniform growth temperature and source flux homogeneity across the 2 in. substrate. Since the substrate directly faces to the evaporation sources, the growth rate is also increased as compared to the lateral growth. The synthesized $MoS_{2-x}O_x$ triangles are spatially uniform as comparing the synthesized $MoS_{2-r}O_r$ triangles with horizontal placed sapphire substrate (Figures S1-S3, Supporting Information). According to the designed reaction routes in Figure 1b,^[25] we use Ar and Ar/O₂ carrier gases for S and MoO₃ sources independently during the growth. Note that O₂ carrier gas could not only balance the sulfurization capacity of S sources but also act as the oxygen doping source for growth of $MoS_{2-x}O_x$.^[18,26,27] Besides, the growth temperature $(T_{\rm G})$, i.e., the temperature of substrates, affects the doping process dominantly.^[25,28] At $T_{\rm G}$ > 900 °C, sulfur has very strong sulfuration properties and MoO₃ can be totally sulfureted to MoS₂ with improved crystal quality, and oxygen doping process is



Figure 1. The growth of monolayer $MoS_{2-x}O_x$ films. a) Schematic set up of a three-temperature zone CVD system for synthesis monolayer $MoS_{2-x}O_x$ films. b) The phase diagram and designed reaction routes for synthesizing $MoS_{2-x}O_x$. c–g) Optical images of the synthesized $MoS_{2-x}O_x$ triangles with increasing doping levels by varying the oxygen carrier gas flow-rate (F_{O2}) from 4 to 25 sccm and growth temperature (T_G) from 930 to 800 °C. Inset in (e) is the AFM image across the boundary of $MoS_{2-x}O_x$ triangle. Scale bar, 1 µm. h) The controllable growth of $MoS_{2-x}O_x$ heterostructures with F_{O2} of 20/6/10/6 sccm at 800 °C.

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Mo4+

S 2s

Mo⁶

a 30

20

10

0

d

ntensity (×1000 a.u.)



Figure 2. The optical properties characterizations of $MoS_{2-x}O_x$. a) XPS spectra of $MoS_{2-x}O_x$ synthesized with $F_{O2} = 6$ sccm. b,c) Raman and PL spectrum of MoS_{2-x}O_x with increased oxygen doping levels. d–f) PL mapping of as-grown MoS_{2-x}O_x samples synthesized with F_{O2} of 4 sccm (d), 10 sccm (e), and 20/6/10/6 sccm (f).

suppressed meanwhile, as comparing optical properties of the $MoS_{2-x}O_x$ samples synthesized with oxygen flow-rate (F_{O2}) from 0 to 20 sccm at 930 °C (Figures S4 and S5, Supporting Information). While the oxygen doping dopants can be reliably introduced and tuned at low T_G , e.g., $T_G < 850$ °C, by just varying F_{O2} (Figures S6 and S7, Supporting Information).

Figure 1c-g shows optical microscope images of typical asgrown MoS_{2-x}O_x triangles obtained under different growth conditions. We can see that their optical contrasts accentuate with F_{O2} , suggesting the formation of $MoS_{2-x}O_x$. Clearer contrast can be found from their undoped outer edges which are grown during cooling down process when O2 gas supply was switched off (Figure 1e). Atomic force microscope (AFM) imaging on these triangles shows no obvious height differences across the $MoS_{2-x}O_x/MoS_2$ boundaries (inset in Figure 1e; Figure S8, Supporting Information), and the 0.7 nm film thickness equals to a monolayer thickness. An excessive oxygen doping could lead to structural damages and etch the $MoS_{2-x}O_x$ basal plane during growth as shown in Figure 1g. This in situ oxygen doping approach can also be used to produce 2D heterostructures in desired locations with controlled dimensions. As shown in Figure 1h, we synthesized lateral heterostructure in the same $MoS_{2-x}O_x$ flake via spatially engineering the oxygen doping levels. The synthesized $MoS_{2-x}O_x$ heterostructure with sharp optical contrast was produced by sequential switching F_{O2} to be 20/6/10/6 sccm during the growth process at 800 °C.

We first performed X-ray photoemission spectroscopy (XPS) for semiqualitative analysis of the oxygen dopant concentrations of $MoS_{2-x}O_x$ samples as shown in Figure 2a and

Figure S9 (Supporting Information). The Mo 3d spectrum is fitted to the doublet peaks at 230 and 233.1 eV with \approx 3.1 eV spin-orbit splitting, corresponding to Mo4+ peaks of Mo-S bonds, the doublet peaks at 230.2 and 233.3 eV, corresponding to Mo⁴⁺ peaks of Mo-O bonds, the doublet peaks at 233.2 and 236.3 eV, corresponding to Mo^{6+} peaks for MoO_3 compounds formed in oxygen-doped MoS_{2-x}O_x samples.^[29,30] Note that the peak at 227.3 eV corresponds to the binding energy of S 2s electrons. Thus, we can qualitatively estimate the amount of oxygen dopant concentrations by considering the peak areas of Mo⁴⁺ of Mo-O bonds and Mo-S bonds, and Mo⁶⁺ of MoO₃ as shown in Table 1. For the O_S concentrations (x), it follows 0%, 0%,16.5%, 22.7%, 24.7%, and 26.1% for $F_{O2} = 0$, 4, 6, 10, 20, and 25 sccm $MoS_{2-x}O_x$ samples according to: S(Mo-O)/[S(Mo-O) + S(Mo-S)] of Mo⁴⁺ species. For the MoO₃ concentrations, it follows 0%, 0%, 5.5%, 10.5%, 17.2%, and 15.8% according to: $S(Mo^{6+})/[S(Mo^{4+}) + S(Mo^{6+})]$.

Table 1. Oxygen dopant concentrations of $MoS_{2-x}O_{x^*}$

F _{O2} [sccm]	T _G [°C]	O _S ratio [%]	MoO ₃ ratio [%]
0	930	0	0
4	930	0	0
6	800	16.5	5.5
10	800	22.7	10.5
20	800	24.7	17.2
25	800	26.1	15.8

0



Then, we further carried out optical spectroscopy characterizations for these monolayer MoS_{2-x}O_x samples with different doping concentrations as illustrated in Figure 1c-g. The Raman spectra are shown in Figure 2b. Beside E_{2g} and A_{1g} modes of intrinsic monolayer MoS₂, a series of new Raman vibration modes appear. The new peaks at 287 cm⁻¹ ($B_{2\sigma}$), 331 cm⁻¹ (B_{1g}), and 464 cm⁻¹ (B_{1g}), corresponding to the vibration modes of Mo–O bonds,^[24,31] confirm the oxygen doping in monolayer MoS₂. According to Hardcastle's empirical relationship^[32] between Mo-O bond length (R, Å) and Raman vibrational frequencies (v, cm⁻¹), which follows: $v(cm^{-1}) =$ $32895\exp(-2.073 \cdot R)$, based on diatomic approximation, we could estimate the Mo-O bond lengths of 2.29 Å from the first observed Raman vibrational frequencies at 287 cm⁻¹, which is close to Mo–S bond length (2.41 Å) in MoS₂ lattice with a \approx 5% shortening. Thus, we assign the Raman peaks at 287 cm⁻¹ to excellent estimates of the vibrational modes of Mo-O bonds for oxygen substitutional atoms in MoS₂ lattice. These new Raman peaks become more prominent with increased F_{O2} , which suggests an increased oxygen dopant concentration and indicates the gradually formation of MoO₃ areas at high oxygen dopant concentrations of $MoS_{2-x}O_x$ samples. Apart from these new modes, increased oxygen doping levels would induce gradual broadening and red/blue shifts of E2g/A1g vibration modes as shown in Figure 2b.^{[29]} PL spectra of ${\rm MoS}_{2-{\rm r}}{\rm O}_{\rm r}$ samples with F_{O2} varying from 0 to 20 sccm were shown in Figure 2c. With the increased oxygen doping levels, the optical bandgap decreases from 1.9 eV at $F_{O2} = 0$ sccm to 1.82 eV at $F_{02} = 10$ sccm, revealing a modulation of the bandgap by oxygen doping. Significant PL quenching is present when F_{02} > 10 sccm. Note that the two prominent sharp PL peaks around 1.78 eV originate from sapphire substrate. Figure 2d-f shows the spatially resolved PL mapping of monolayer $MoS_{2-x}O_x$ triangles with F_{O2} of 4 and 10 sccm, as well as $MoS_{2-x}O_x$ heterostructure with F_{O2} of 20/6/10/6 sccm. The good uniformity of PL intensity in Figure 2d,e indicates the spatially homogeneous doping of the entire $MoS_{2-x}O_x$ triangles. Besides, the sharp PL intensity contrast of the synthesized $MoS_{2-x}O_x$ heterostructures (Figure 2f) indicates the successful synthesizing of in-plane $MoS_{2-x}O_x$ heterostructures with sharp boundaries. Thus, we are able to engineer the electronic structure of $MoS_{2-x}O_x$ by spatial modulation the oxygen dopant concentrations.

In order to reveal the atomic structure and local electronic properties of monolayer $MoS_{2-x}O_x$ samples, we thus performed STM/STS and STEM experiments. Monolayer MoS_{2-x}O_x film was grown on the highly oriented pyrolytic graphite (HOPG) with $F_{O2} = 6$ sccm at 800 °C as shown in **Figure 3**a. Figure 3b,c shows typical atomic resolution STM images of monolayer $MoS_{2-r}O_r$ at different scanning scales. The uniformly distributed hexagonal halo patterns in Figure 3b are assigned to Os substitutional sites in the MoS2 crystal lattice, in agreement with previous reports.^[15,16,18] Importantly, two different halo patterns having obvious contrasts than the background MoS₂ lattice are visualized. Our simulated STM images based on density-functional theory (DFT) calculations (details in the Experimental Section) show that these two features, i.e., the darker/brighter, correspond to substitutional O atom in the top/ bottom S layer (O_{S-top}/O_{S-bottom}), as illustrated in the suggested atomic structures and simulated STM images in Figure 3d,e.



Figure 3. Atomic structure configuration and localized electronic properties of O_S in MoS_{2-x}O_x. a) AFM image of as-grown MoS_{2-x}O_x on HOPG substrates (F_{O2} = 6 sccm, 800 °C). b,c) High-resolution STM images of MoS_{2-x}O_x at different scale. Set points: 0.74 V, 100 pA (b); 0.8 V, 50 pA (c). The O_S could be distinguished by dark (O_{S-top}) or bright (O_{S-bottom}) hexagonal halo patterns. d,e) Simulated STM images of O_{S-top} (d) and O_{S-bottom} (e). Left panels correspond to top and side view of the DFT-calculated atomic structure and O_S sites are marked by black circles. f) Representative dI/dV spectra acquired on O_S of MoS_{2-x}O_x and pristine MoS₂. Set points: 1.1 V, 50 pA.

From our STM images, it shows clearly that the substitutional sites are predominately on the top S layer of $MoS_{2-x}O_x$. The O-atom density is estimated to be in the range of 10¹³–10¹⁴ cm⁻². Compared to the pristine MoS₂, the introduction of O_S atoms into it clearly modifies its local electronic structure, evidenced from our STS measurements. Figure 3f shows the dI/dV spectra taken on O_S sites in monolayer $MoS_{2,x}O_x$ together with a control spectrum from pristine MoS₂ (Figure S10, Supporting Information). The measured electronic bandgap of pristine MoS₂ is ≈2.25 eV and the bandgap reduced to ≈1.75 eV of $MoS_{2-x}O_x$ sample (F₀₂ = 6 sccm). In general, sulfur vacancies would introduce in-gap states in STS spectra;[33,34] while there are no obvious in-gap states observed in the measured STS spectra, suggesting that oxygen dopants are in a form of substitution chalcogen atoms in our CVD grown monolayer $MoS_{2-x}O_x$ samples and this substitutional doping can effectively modulate their bandgaps. Further, we carried out atomic structure characterizations of as-grown $MoS_{2-x}O_x$ samples with different doping levels as shown in Figure 4. Though it is difficult to distinguish the light substitutional atoms of O_S (S–O) and S vacancy (S–)



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Figure 4. The atomic structural integrity of $MoS_{2-x}O_x$ with increased oxygen doping levels. a–d) As-recorded high resolution ADF-STEM images and e–h) SAED patterns of $MoS_{2-x}O_x$ samples.

due to their relatively weak intensity contrast (Figures S12 and S13, Supporting Information),^[35] the overall structural integrity and periodical atomic structure of 2H phase still maintain even at high doping levels as confirmed by high-resolution STEM images (Figure 4a–d; Figure S14, Supporting Information) and selected area electron diffraction (SAED) pattern (Figure 4e–h), which indicate that our doping strategy is nondestructive technique for realizing controllable doping in MoS₂ crystal lattice.

As mentioned above, the oxygen substitutional doping in MoS₂ can significantly affect their electronic properties. We also fabricated FETs based on these MoS_{2-x}O_x monolayers to study their electronic transport properties as shown in Figure 5a. Refer to the Experimental Section and Figure S15 (Supporting Information) for details on the device structure and fabrication. Mechanical exfoliated few-layer graphene (FLG) was used as the source/drain and gate electrodes since graphene contact^[36,37] has low contact resistance (R_c) of 2.96 k Ω µm as confirmed in Figure S16 of the Supporting Information. Thin hBN flakes with atomically smooth surfaces and free of dangling bonds were used as the dielectric layers to effectively avoid impurity scatterings and trapped charges.^[38] Figure 5b,c shows output $(I_{ds}-V_{ds})$ and transfer $(I_{ds}-V_{g})$ curves, respectively, of a typical $MoS_{2-x}O_x$ device (F₀₂ = 6 sccm). The on-current density (I_{ds}/W) is as high as $\approx 90 \ \mu A \ \mu m^{-1}$. The linear Ids-Vds characteristics at small biases suggest an excellent contact between $MoS_{2-x}O_x$ and FLG. The saturation characteristics of Ids-Vds at large biases and wide range tunable threshold voltages are important features when interconnecting them for integrated circuits. High on/off ratio of \approx 4.5 \times 10⁹, low off-current \approx 100 fA, and ideal subthreshold swing (SS) of 80 mV dec⁻¹ can also be achieved in these devices, indicating outstanding device performances. Figure 5d shows the transfer curves of several $MoS_{2-x}O_x$ FETs with $F_{O2} = 4$, 6, 10, and 20 sccm. An obvious shift of subthreshold voltages $(V_{\rm th})$ by doping (n-type) levels can be clearly seen. We also calculated the field effect mobility ($\mu_{\rm FF}$) of these MoS_{2-x}O_x FETs at different doping levels (Figure 5e). Surprisingly, μ_{FE} can be significantly improved by moderate doping, e.g., from 30 cm² V⁻¹ s⁻¹ at $F_{O2} = 4$ sccm (which is close to an intrinsic case), to 70 cm² V⁻¹ s^{-1} at $F_{O2} = 6$ sccm, and a maximum value of 105 cm² V⁻¹ s⁻¹ can be achieved at $F_{O2} = 10$ sccm (Figures S17– S21, Supporting Information, for detailed electrical analysis). As the highly conductive phase of $MoS_{2-x}O_x$ from O_S substitutional doping could be formed into MoS_2 lattice,^[9,14] the electronic behavior of doped MoS₂ films could be modulated, and the sheet conductance and associated field effect mobility $(\mu_{\rm FE})$ could be enhanced with appropriate doping levels (F_{Q2} \leq 10 sccm) as shown in Figure 5e and Figure S22 (Supporting Information). At heavy doping levels ($F_{O2} \ge 20$ sccm), the defects and visible damages (Figure 1g) prone to form on the $MoS_{2-x}O_x$ basal plane and serve as the main carrier scattering centers, thus degrading its electronic quality. Indeed, an abrupt dropping of μ_{FE} when $F_{O2} > 20$ sccm was observed. Furthermore, SS of devices also increases gradually due to the enhanced oxygen doping concentration (Figure 5d). As highlighted in Figure 5f, these moderate doped $MoS_{2-x}O_x$ monolayers show much higher mobilities and on/off ratios than the intrinsic monolayer or few layer MoS_2 .^[3,39–42]

Based on these high-quality $MoS_{2-x}O_x$ monolayers, we also fabricated complexed logic gates. **Figure 6**a–d shows as-fabricated inverters, NAND, NOR, and AND gates by integrating 2, 3, 3, and 5 FETs, respectively. The output voltage of an inverter has a sharp switching characteristic when sweeping the input voltage (Figure 6e). The voltage gain is 90 at $V_{dd} = 5$ V, which is superior to pristine MoS_2 .^[43] Logic functions of NAND, NOR, ADVANCED SCIENCE NEWS _____



Figure 5. Electrical performances of $MoS_{2-x}O_x$ FETs. a) Cross-section schematic illustration and optical image of $MoS_{2-x}O_x$ FETs with few-layer graphene (FLG) contact. b,c) Output and transfer curves of a $MoS_{2-x}O_x$ device with $F_{O2} = 6$ sccm. Channel length/width (*L/W*) is 5/5 μ m and hBN thickness (*t*) is 31 nm. d) Normalized transfer curves with 1 V bias as a function of carrier density ($n = C_i V_g$) of $MoS_{2-x}O_x$ FETs with different oxygen doping levels. Solid/dashed lines correspond to left/right panels in logarithm/linear scales. e) Device mobility (μ_{FE}) of $MoS_{2-x}O_x$ FETs at different oxygen doping levels. f) Comparisons of μ_{FE} and on/off ratio with literature values.



Figure 6. Demonstration of $MoS_{2-x}O_x$ logic gates devices. a–d) Optical images of (a) inverter, (b) NAND, (c) NOR, (d) AND gates with $MoS_{2-x}O_x$ ($F_{O2} = 6$ sccm) as channel materials. Insets are corresponding schematics of logic gates. e) Output voltage V_{out} (left axis) and voltage gain (right axis) as a function of input voltages V_{in} with $V_{dd} = 5$ V. f–h) Output voltages of the logic NAND (f), NOR (g), and AND (h) gates at four typical input states with $V_{dd} = 2$ V. "0" and "1" labels represent the low and high binary output states, individually.

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and AND could be reliably realized at four typical input states ($V_A = 0$ (0 V) or 1 (2 V), $V_B = 0$ or 1) as shown in Figure 6f–h. These results suggest that these MoS_{2-x}O_x monolayers with tunable electronic properties could serve as promising 2D building blocks for novel electronics with better performances than that of the intrinsic MoS₂ monolayers.

3. Conclusion

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In summary, we developed an effective CVD technique for in situ oxygen doping of monolayer MoS_2 in a controlled manner and oxygen dopants mainly adopt a form of substitution chalcogen atoms in monolayer $MoS_{2-x}O_x$ samples. We found that the oxygen doping leads to a significant modulation on their optical and electronic properties. Greatly improved carrier mobilities and controllable n-type electronic doping effects resulted from optimized oxygen doping levels of $MoS_{2-x}O_x$ were revealed. The $MoS_{2-x}O_x$ -based devices behave excellent electronic performances comparing with intrinsic MoS_2 monolayers, suggesting that such high-quality monolayer films could be better alternatives for future electronic applications.

4. Experimental Section

Growth of $MoS_{2-x}O_x$: The growth of monolayer $MoS_{2-x}O_x$ was carried out in a three-temperature zone CVD system (Figure 1a). S powder (Alfa, 99.5%, 8 g) and MoO_3 powder (Alfa, 99.9995%, 30 mg) were loaded in two inner quartz tube before growth and placed at zone I and zone II, and c-plane polished sapphire (or newly exfoliated graphite) was placed at zone III as substrate. Note that sapphire substrates were preannealed in O_2 atmosphere at 1000 °C for 4 h to form atomic flat surfaces. The carrier gas of Ar (40 sccm) and Ar/O_2 (240/0–25 sccm) was fluxed for S power and MoO_3 sources individually. The pressure in the chamber was kept at \approx 1 Torr. Temperatures at zone I, zone II, and zone III were 130, 530, and 800–930 °C and maintained for 30–50 min during the growth.

Sample Characterizations and Measurements: Raman and PL spectra were carried out by a micro-Raman spectrometer (Horiba LabRAM HR Evolution) in a confocal backscattering configuration with an excitation laser wavelength of 532 nm, a laser power of 1 mW and spot size of $1\ \mu\text{m}^2.$ The PL mapping was obtained with WiTec (Alpha 300R) Raman microscope with a laser power of 1 mW and step size of 300 nm. The XPS was acquired from Kratos Analytical Axis Ultra system. STM experiments were performed in a homebuilt low temperature-STM system (4.5 K, 10⁻¹⁰ Torr) with chemical etched W tip. STM data were processed by using the free WSxM software.[44] The STS spectra were taken by using a lock-in amplifier with a modulation voltage of 20 mV and at a frequency of 677 Hz. STEM characterizations were acquired by an aberration-corrected Nion U-HERMES200 system operated at 60 kV and STEM simulation was carried out by COMPUTEM software. AFM measurements were performed by Asylum Research Cypher S instrument using RTESP-300 tips. All the electrical measurements were carried out in a lanus probe station (base pressure $\approx 10^{-6}$ Torr) with Agilent semiconductor parameter analyzer (4156C and B1500, high resolution modules) at room temperature.

DFT Calculations: First-principles calculations were performed for top and bottom O_S atoms in the MoS₂-basal plane by using the Quantum-ESPRESSO software package,^[45,46] which is based on the DFT. The ultrasoft pseudopotential^[47] and exchange-correlation functional within a generalized gradient approximation of Perdew–Burke–Ernzerh formula have been used.^[48] The sampling of Brillouin zone with a 16 × 16 × 1 k-point mesh was used for a 5 × 5 supercell with a vacuum layer of ≈15 Å for simulations.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

Acknowledgements

J.T. and Z.W. contributed equally to this work. G.Y.Z. and R.Y. supervised this project. Z.W., Q.Q.W., and B.Y.H. carried out the growth of MoS_{2-x}O_x samples. Y.W. and K.H.W. performed STM characterizations. J.T., B.H., X.M.L., and P.G. prepared the STEM samples and performed characterizations, analysis and simulations. J.T. performed the optical spectroscopy measurements, fabricated transistors and logic gates, and carried out electrical measurements. J.T., R.Y., and G.Y.Z. analyzed data, wrote the manuscript and all authors discussed and commented on it. The authors thank Kenii Watanabe and Takashi Taniguchi in NIMS (Japan) for providing high quality BN samples, Xuanyi Li, Jiatao Sun, Zhen Chi, Hailong Chen, Wu Zhou, and Lin Gu in IOP-CAS (China) for fruitful discussions. The authors acknowledge the supports from the National Science Foundation of China (NSFC, Grant Nos. 61888102, 11834017, 61734001, 51572289, 11825405, 1192780039, 51672007, and 11974023), the Strategic Priority Research Program of Chinese Academy of Sciences (CAS, Grant No. XDB0000000), the Key Research Program of Frontier Sciences of CAS (Grant No. QYZDB-SSW-SLH004), the National Key R&D program (Grant No. 2016YFA0300904), the Research Program of Beijing Academy of Quantum Information Sciences (Grant No. Y18G11), and the Youth Innovation Promotion Association of CAS (Grant No. 2018013).

Conflict of Interest

The authors declare no conflict of interest.

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2D electronics, bandgap engineering, MoS₂, substitutional doping

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- G. Fiori, F. Bonaccorso, G. Iannaccone, T. Palacios, D. Neumaier, A. Seabaugh, S. K. Banerjee, L. Colombo, *Nat. Nanotechnol.* 2014, 9, 768.
- [2] Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman, M. S. Strano, Nat. Nanotechnol. 2012, 7, 699.
- [3] K. Kang, S. Xie, L. Huang, Y. Han, P. Y. Huang, K. F. Mak, C. J. Kim, D. Muller, J. Park, *Nature* 2015, 520, 656.
- [4] Z. Lin, Y. Liu, U. Halim, M. Ding, Y. Liu, Y. Wang, C. Jia, P. Chen, X. Duan, C. Wang, F. Song, M. Li, C. Wan, Y. Huang, X. Duan, *Nature* **2018**, *562*, 254.
- [5] D. Rhodes, S. H. Chae, R. Ribeiro-Palau, J. Hone, Nat. Mater. 2019, 18, 541.
- [6] Y. Gong, H. Yuan, C. L. Wu, P. Tang, S. Z. Yang, A. Yang, G. Li, B. Liu, J. van de Groep, M. L. Brongersma, M. F. Chisholm, S. C. Zhang, W. Zhou, Y. Cui, *Nat. Nanotechnol.* **2018**, *13*, 294.
- [7] S. Z. Yang, Y. Gong, P. Manchanda, Y. Y. Zhang, G. Ye, S. Chen, L. Song, S. T. Pantelides, P. M. Ajayan, M. F. Chisholm, W. Zhou, *Adv. Mater.* **2018**, *30*, 1803477.
- [8] H. Li, X. Duan, X. Wu, X. Zhuang, H. Zhou, Q. Zhang, X. Zhu,
 W. Hu, P. Ren, P. Guo, L. Ma, X. Fan, X. Wang, J. Xu, A. Pan,
 X. Duan, J. Am. Chem. Soc. 2014, 136, 3756.

ADVANCED SCIENCE NEWS

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- [9] J. Jadwiszczak, C. O'Callaghan, Y. Zhou, D. S. Fox, E. Weitz, D. Keane, C. P. Cullen, I. O'Reilly, C. Downing, A. Shmeliov, P. Maguire, J. J. Gough, C. McGuinness, M. S. Ferreira, A. L. Bradley, J. J. Boland, G. S. Duesberg, V. Nicolosi, H. Zhang, *Sci. Adv.* **2018**, *4*, 5031.
- [10] F. Zhang, Y. Lu, D. S. Schulman, T. Zhang, K. Fujisawa, Z. Lin, Y. Lei, A. L. Elias, S. Das, S. B. Sinnott, M. Terrones, *Sci. Adv.* **2019**, *5*, 5003.
- [11] D. Kiriya, M. Tosun, P. Zhao, J. S. Kang, A. Javey, J. Am. Chem. Soc. 2014, 136, 7853.
- [12] L. Gao, Q. Liao, X. Zhang, X. Liu, L. Gu, B. Liu, J. Du, Y. Ou, J. Xiao, Z. Kang, Z. Zhang, Y. Zhang, Adv. Mater. 2020, 32, 1906646.
- [13] P. D. Fleischauer, J. R. Lince, Tribol. Int. 1999, 32, 627.
- [14] H. Gao, J. Suh, M. C. Cao, A. Y. Joe, F. Mujid, K. H. Lee, S. Xie, P. Poddar, J. U. Lee, K. Kang, P. Kim, D. A. Muller, J. Park, *Nano Lett.* 2020, 20, 4095.
- [15] S. Barja, S. Refaely-Abramson, B. Schuler, D. Y. Qiu, A. Pulkin, S. Wickenburg, H. Ryu, M. M. Ugeda, C. Kastl, C. Chen, C. Hwang, A. Schwartzberg, S. Aloni, S.-K. Mo, D. Frank Ogletree, M. F. Crommie, O. V. Yazyev, S. G. Louie, J. B. Neaton, A. Weber-Bargioni, *Nat. Commun.* **2019**, *10*, 3382.
- [16] B. Schuler, J. H. Lee, C. Kastl, K. A. Cochrane, C. T. Chen, S. Refaely-Abramson, S. Yuan, E. van Veen, R. Roldan, N. J. Borys, R. J. Koch, S. Aloni, A. M. Schwartzberg, D. F. Ogletree, J. B. Neaton, A. Weber-Bargioni, ACS Nano 2019, 13, 10520.
- [17] A. W. Robertson, Y. C. Lin, S. Wang, H. Sawada, C. S. Allen, Q. Chen, S. Lee, G. D. Lee, J. Lee, S. Han, E. Yoon, A. I. Kirkland, H. Kim, K. Suenaga, J. H. Warner, ACS Nano 2016, 10, 10227.
- [18] J. Peto, T. Ollar, P. Vancso, Z. I. Popov, G. Z. Magda, G. Dobrik, C. Hwang, P. B. Sorokin, L. Tapaszto, *Nat. Chem.* 2018, *10*, 1246.
- [19] S. Susarla, A. Kutana, J. A. Hachtel, V. Kochat, A. Apte, R. Vajtai, J. C. Idrobo, B. I. Yakobson, C. S. Tiwary, P. M. Ajayan, *Adv. Mater.* 2017, *29*, 1702457.
- [20] Y. Gong, Z. Liu, A. R. Lupini, G. Shi, J. Lin, S. Najmaei, Z. Lin, A. L. Elías, A. Berkdemir, G. You, H. Terrones, M. Terrones, R. Vajtai, S. T. Pantelides, S. J. Pennycook, J. Lou, W. Zhou, P. M. Ajayan, *Nano Lett.* **2014**, *14*, 442.
- [21] J. Zhou, J. Lin, X. Huang, Y. Zhou, Y. Chen, J. Xia, H. Wang, Y. Xie, H. Yu, J. Lei, D. Wu, F. Liu, Q. Fu, Q. Zeng, C. H. Hsu, C. Yang, L. Lu, T. Yu, Z. Shen, H. Lin, B. I. Yakobson, Q. Liu, K. Suenaga, G. Liu, Z. Liu, *Nature* **2018**, *556*, 355.
- [22] N. Kang, H. P. Paudel, M. N. Leuenberger, L. Tetard, S. I. Khondaker, J. Phys. Chem. C 2014, 118, 21258.
- [23] S. I. Khondaker, M. R. Islam, J. Phys. Chem. C 2016, 120, 13801.
- [24] J. Xie, J. Zhang, S. Li, F. Grote, X. Zhang, H. Zhang, R. Wang, Y. Lei, B. Pan, Y. Xie, J. Am. Chem. Soc. 2013, 135, 17881.
- [25] Q. Ji, Y. Zhang, Y. Zhang, Z. Liu, Chem. Soc. Rev. 2015, 44, 2587.
- [26] W. Chen, J. Zhao, J. Zhang, L. Gu, Z. Yang, X. Li, H. Yu, X. Zhu, R. Yang, D. Shi, X. Lin, J. Guo, X. Bai, G. Zhang, J. Am. Chem. Soc. 2015, 137, 15632.
- [27] H. Yu, M. Liao, W. Zhao, G. Liu, X. J. Zhou, Z. Wei, X. Xu, K. Liu, Z. Hu, K. Deng, S. Zhou, J. A. Shi, L. Gu, C. Shen, T. Zhang, L. Du, L. Xie, J. Zhu, W. Chen, R. Yang, D. Shi, G. Zhang, ACS Nano 2017, 11, 12001.
- [28] K. Liu, W. Zhang, Y. Lee, Y. Lin, M. Chang, C. Su, C. Chang, H. Li, Y. Shi, H. Zhang, C. Lai, L. Li, *Nano Lett.* **2012**, *12*, 1538.
- [29] T. Y. Ko, A. Jeong, W. Kim, J. Lee, Y. Kim, J. E. Lee, G. H. Ryu, K. Park, D. Kim, Z. Lee, M. H. Lee, C. Lee, S. Ryu, 2D Mater. 2017, 4, 014003.
- [30] M. A. Baker, R. Gilmore, C. Lenardi, W. Gissler, Appl. Surf. Sci. 1999, 150, 255.

[31] L. Seguin, M. Figlarz, R. Cavagnat, J.-C. Lassegues, Spectrochim. Acta 1995, 51, 1323.

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- [32] F. D. Hardcastle, I. E. Wachs, J. Raman. Spectrosc. 1990, 21, 683.
- [33] B. Schuler, D. Y. Qiu, S. Refaely-Abramson, C. Kastl, C. T. Chen, S. Barja, R. J. Koch, D. F. Ogletree, S. Aloni, A. M. Schwartzberg, J. B. Neaton, S. G. Louie, A. Weber-Bargioni, *Phys. Rev. Lett.* **2019**, 123, 076801.
- [34] P. Vancso, G. Z. Magda, J. Peto, J. Y. Noh, Y. S. Kim, C. Hwang, L. P. Biro, L. Tapaszto, *Sci. Rep.* **2016**, *6*, 29726.
- [35] O. L. Krivanek, M. F. Chisholm, V. Nicolosi, T. J. Pennycook, G. J. Corbin, N. Dellby, M. F. Murfitt, C. S. Own, Z. S. Szilagyi, M. P. Oxley, S. T. Pantelides, S. J. Pennycook, *Nature* 2010, 464, 571.
- [36] L. Xie, M. Liao, S. Wang, H. Yu, L. Du, J. Tang, J. Zhao, J. Zhang, P. Chen, X. Lu, G. Wang, G. Xie, R. Yang, D. Shi, G. Zhang, *Adv. Mater.* **2017**, *29*, 1702522.
- [37] X. Cui, G. H. Lee, Y. D. Kim, G. Arefe, P. Y. Huang, C. H. Lee, D. A. Chenet, X. Zhang, L. Wang, F. Ye, F. Pizzocchero, B. S. Jessen, K. Watanabe, T. Taniguchi, D. A. Muller, T. Low, P. Kim, J. Hone, *Nat. Nanotechnol.* **2015**, *10*, 534.
- [38] C. R. Dean, A. F. Young, I. Meric, C. Lee, L. Wang, S. Sorgenfrei, K. Watanabe, T. Taniguchi, P. Kim, K. L. Shepard, J. Hone, *Nat. Nanotechnol.* 2010, *5*, 722.
- [39] B. Radisavljevic, A. Kis, Nat. Mater. 2013, 12, 815.
- [40] S. L. Li, K. Wakabayashi, Y. Xu, S. Nakaharai, K. Komatsu, W. W. Li, Y. F. Lin, A. Aparecido-Ferreira, K. Tsukagoshi, *Nano Lett.* 2013, *13*, 3546.
- [41] R. Kappera, D. Voiry, S. E. Yalcin, B. Branch, G. Gupta, A. D. Mohite, M. Chhowalla, *Nat. Mater.* **2014**, *13*, 1128.
- [42] S. Kim, A. Konar, W. S. Hwang, J. H. Lee, J. Lee, J. Yang, C. Jung, H. Kim, J. B. Yoo, J. Y. Choi, Y. W. Jin, S. Y. Lee, D. Jena, W. Choi, K. Kim, *Nat. Commun.* **2012**, *3*, 1011.
- [43] H. Wang, L. Yu, Y. H. Lee, W. Fang, A. Hsu, P. Herring, M. Chin, M. Dubey, L. J. Li, J. Kong, T. Palacios, *Int. Electron Devices Meeting*, 2012 (IEDM 2012), San Francisco, CA, USA, IEEE, CA 2012, p. 4.6.1.
- [44] I. Horcas, R. Fernandez, J. M. Gomez-Rodriguez, J. Colchero, J. Gomez-Herrero, A. M. Baro, *Rev. Sci. Instrum.* 2007, 78, 013705.
- [45] P. Giannozzi, S. Baroni, Bonini, N. M. Calandra, R. Car, C. Cavazzoni, D. Ceresoli, G. L. Chiarotti, M. Cococcioni, I. Dabo, A. Dal Corso, S. de Gironcoli, S. Fabris, G. Fratesi, R. Gebauer, U. Gerstmann, C. Gougoussis, A. Kokalj, M. Lazzeri, L. Martin-Samos, N. Marzari, F. Mauri, R. Mazzarello, S. Paolini, A. Pasquarello, L. Paulatto, C. Sbraccia, S. Scandolo, G. Sclauzero, A. P. Seitsonen, A. Smogunov, P. Umari, R. M. Wentzcovitch, J. Phys.: Condens. Matter 2009, 21, 395502.
- [46] P. Giannozzi, O. Andreussi, T. Brumme, O. Bunau, M. B. Nardelli, M. Calandra, R. Car, C. Cavazzoni, D. Ceresoli, M. Cococcioni, N. Colonna, I. Carnimeo, A. Dal Corso, S. de Gironcoli, P. Delugas, R. A. DiStasio, A. Ferretti, A. Floris, G. Fratesi, G. Fugallo, R. Gebauer, U. Gerstmann, F. Giustino, T. Gorni, J. Jia, M. Kawamura, H. Y. Ko, A. Kokalj, E. Kucukbenli, M. Lazzeri, M. Marsili, N. Marzari, F. Mauri, N. L. Nguyen, H. V. Nguyen, A. Otero-de-la-Roza, L. Paulatto, S. Ponce, D. Rocca, R. Sabatini, B. Santra, M. Schlipf, A. P. Seitsonen, A. Smogunov, I. Timrov, T. Thonhauser, P. Umari, N. Vast, X. Wu, S. Baroni, J. Phys.: Condens. Matter 2017, 29, 465901.
- [47] D. Vanderbilt, Phys. Rev. B: Condens. Matter Mater. Phys. 1990, 41, 7892.
- [48] J. P. Perdew, K. Burke, M. Ernzerhof, Phys. Rev. Lett. 1996, 77, 3865.